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Thermal properties and devitrification behavior of $(2.5-x)CaO \cdot x/3La_2O_3 \cdot 2SiO_2$

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Abstract

La₂O₃ was substituted for CaO in a glass of composition 2.5CaO·2SiO₂ and the effect of this substitution on thermal properties and non-isothermal behavior was studied. The glass samples were submitted to thermal analysis (DTA) and X-ray diffraction analysis. The trends with the composition of the glass transformation, T_g , and the activation energy of crystal growth, E_c , are the result of a type of competition between Ca²⁺ and La³⁺ ions each to achieve their preferred coordination. The lanthanum ion does not succeed in realizing its own average coordination. As a result $T_g=810\pm5^{\circ}$ C and the activation energy for crystal growth was $E_c=650\pm25$ kJ/mol, regardless of composition. Devitrification involved a mechanism of surface nucleation: however, because the samples softened and sintered before devitrifying, surface nuclei behaved as bulk nuclei. In all the devitrified samples CaO·SiO₂ crystals form together with a secondary calcium silicate phase, while Ca₃La₆(SiO₄)₆ crystals form when \cong 40% molar weight CaO is substituted. The pattern for the La₂O₃ free glass shows that the α CaO·SiO₂ is also formed in the temperature range 900–1000°C, which should be stable only above 1125°C as reported in the phase diagram. © 2000 Elsevier Science B.V. All rights reserved.

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1. Introduction

It is known that glasses of the system CaO-SiO₂ with a CaO molar fraction in the range 0.3–0.55 are bioactive, that is to say they bond to living bone through a thin layer rich in calcium and phosphorus [1,2]. The addition of La₂O₃ to silicates improves both elastic modulus and hardness of the glasses [3]. Recently, it was found [4,5] that the non-isothermal devitrification of glasses of composition CaO-La₂O₃(MgO)-SiO₂ exhibited peculiar characteristics. These devitrify through a surface nucleation mechanism. However, since in the temperature range of

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efficient crystal growth, softening and sintering occur, surface nuclei behave virtually as bulk nuclei. This was well supported by SEM observations for samples devitrified during a DSC run [6].

In this paper the thermal properties and the study of the devitrification behavior of glasses of composition (2.5-x)CaO·x/3La₂O₃·2SiO₂ are reported. The substitution kept the O/Si molar ratio constant.

2. Experimental

Glasses of composition expressed by the following formula

$$(2.5 - x)$$
CaO $\cdot x/3$ La₂O₃ $\cdot 2$ SiO₂

were prepared by melting analytical grade reagents, La_2O_3 , $CaCO_3$ and SiO_2 in a platinum crucible in an

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electric furnace for 4 h, in the temperature range $1400-1600^{\circ}$ C. The melts were quenched by plunging the bottom of the crucible into cold water. The products were weighted after quenching and it was found that losses were not greater than 0.5%.

Differential thermal analysis (DTA) was carried out by means of a Netzsch differential scanning calorimeter (DSC) model 404M on approximately 50 mg powdered samples at various heating rates (2–20°C/ min). Finely (63–90 μ m) and coarsely (315–500 μ m) powdered samples were used. Powdered Al₂O₃ was used as reference material. In order to calibrate the instrument, a preliminary DTA was carried out using Al₂O₃. The DTA curve obtained did not show any inflection in the baseline.

Devitrified samples were analyzed by computerinterfaced X-Ray (Cu K α) powder diffractometry (XRD) using a Philips Diffractometer model PW1710, with a scan speed of 1°/min and a built-in computer search program. The crystalline phases were identified by means of JCPDS cards.

3. Results

Figs. 1 and 2 show the DTA curves obtained for finely and coarsely powdered samples, respectively. As illustrated in Fig. 3, the glass transformation temperature, T_g , was taken as the maximum of the DDTA curve in the glass transformation range. In the same figure, the DTA exothermic peak temperature T_p



Fig. 2. DTA curves recorded at 10° C/min on coarsely (315–500 μ m) powdered samples.

is also shown. The uncertainty of both values is of approximately $\pm 1^{\circ}$ C. In Fig. 4 T_{g} and T_{p} values are plotted against the glass composition expressed as the *x* values of the general formula of the series. The substitution appears to have no effect on T_{g} , while a slight increasing trend was observed for T_{p} .

Fig. 5 shows the X-ray diffraction patterns of the samples submitted to a DTA run stopped just after the exothermic peak. The lines were attributed by means of the JCPDS cards. In all patterns, the lines of wollastonite (JCPDS card 27/88) appear. In the (a) pattern (x=0), the lines of pseudo-wollastonite, α CaO·SiO₂ (JCPDS card 19/248) and 3CaO·SiO₂ (JCPDS card 11-593) are also present. It is interesting



Fig. 1. DTA curves recorded at 10°C/min on finely (63-90 µm) powdered samples.



Fig. 3. A typical DTA curve and its first derivative (DDTA).

to observe that α CaO·SiO₂ is the high temperature phase, in fact in the phase diagram CaO·SiO₂ is reported to be stable above 1125°C. In the ((b), (c) and (d)) patterns, they are partially substituted by the lines of 2CaO·SiO₂ (JCPDS card 24-234); in the (c) and (d) patterns (*x*=0.6, 1.0) Ca₃La₆(SiO₄)₆ crystal (JCPDS card 27/78) reflections appear.

Non-isothermal devitrification was also studied. The kinetic parameters were determined using the following two equations:

$$\ln\beta = -\frac{E_{\rm c}}{RT_{\rm p}} + \text{const} \tag{1}$$



Fig. 4. Glass transformation temperature, T_g ($\mathbf{\nabla}$) and peak temperature T_p for finely ($\mathbf{\square}$) and coarsely powdered ($\mathbf{\odot}$) samples vs. composition (*x*).

$$\ln\Delta T = -\frac{mE_c}{RT} + \text{const}$$
(2)

that can be derived from the well-known following expression [7,8]:

$$-\ln(1-\alpha) = \left(\frac{AN}{\beta^m}\right) \exp\left(-\frac{mE_c}{RT}\right)$$
(3)

where α is the crystallization degree, *N* is the nuclei number, *A* is a constant, β is the heating rate, ΔT and T_p are, respectively, the deflection from the baseline and the peak temperature taken as indicated in Fig. 3.



Fig. 5. X-ray diffraction patterns of samples devitrified during a DTA run: (\bigcirc) wollastonite; (\bigcirc) pseudo-wollastonite; (\blacksquare) 3CaOSiO₂; (\square) 2CaOSiO₂; (\diamondsuit) Ca₃La₆ (SiO₄)₆.

Since in silicate glasses the devitrification exo-peak occurs in a higher temperature range than that of efficient nucleation [4], E_c is the crystal growth activation energy. The parameter *m* depends on the mechanism and morphology of crystal growth; it ranges from m=1 for one-dimensional growth to m=3 for three-dimensional growth [7,8].

Eqs. (1) and (2) can be derived from Eq. (3) by supposing that: (1) the value of α at peak temperature is not dependent on the heating rate [9]; (2) ΔT is proportional to the instantaneous reaction rate [10,11]; (3) in the initial part of the DTA crystallization peak the deflection from the baseline, ΔT , is assumed to be constant with the crystallization degree, α , since the change in the temperature has a much greater effect on it [12]. Eq. (1) was obtained using Kissinger method. It was compared to Ozawa method [13,14] and it was shown that, since the changes of β are dominant with respect to $T_{\alpha 0}^{2n}$, equation

$$\frac{1}{n}\ln\left(\frac{\beta^{n+1}}{T_{\alpha 0}^{2n}}\right) = c2 - \frac{E}{R}\frac{1}{T_{\alpha 0}}$$
(4)

can be reduced to

$$\frac{1}{n}\ln\beta^{n+1} = c2 - \frac{E}{R}\frac{1}{T_{x0}}$$
(5)

In Fig. 6 the plots of $\ln \beta$ versus $1/T_p$ are reported. In Figs. 7 and 8 the plots of $\ln \Delta T$ versus 1/T are reported for finely and coarsely powdered samples. According to Eqs. (1) and (2) straight lines were obtained. Their slopes allow the determination of values of E_c and mE_c



Fig. 6. Kissinger plot of $\ln \beta$ vs. $1/T_p$.



Fig. 7. ln ΔT vs. 1/T for finely powdered (63–90 µm) samples.



Fig. 8. $\ln \Delta T$ vs. 1/T for coarsely powdered (315–500 µm) samples.



Fig. 9. Activation energy of crystal growth, E_c vs. composition (x).



Fig. 10. mE_c values for finely (63–90 µm) (\bullet) and coarsely powdered (315–500 µm) (\blacksquare) samples vs. composition (*x*).



Fig. 11. *m* values for finely $(63-90 \,\mu\text{m})$ (\blacksquare) and coarsely powdered $(315-500 \,\mu\text{m})$ (\bigcirc) samples vs. composition (*x*).

with an error of $\pm 10\%$ [14] and, therefore, the magnitudes of *m* for each glass. In Figs. 9–11, E_c , mE_c and *m* values, respectively, are reported as a function of the composition. Only slight variations of E_c values are obtained as *x* increases. It is interesting to observe that the *m* values relative to finely powdered samples are greater than the ones relative to coarsely powdered samples.

4. Discussion

The role of SiO_2 and CaO in the glass structure is well known: the former is a network former oxide,

while the latter in a network modifier. La₂O₃ is a less common component of glasses. The criterion reported in the literature [18,20] suggest that it is expected to be a network modifier. This role proved to be useful in explaining T_g versus composition curves of glasses of the system Na₂O-M₂O₃-SiO₂ (M=La, Sc, Y, Al, B) obtained by substituting M₂O₃ for Na₂O starting from the glass composition Na₂O·2SiO₂ [19].

Following Ray [17], T_g values depend on the density of covalent cross-linking and the numbers and strengths of cross-linking between the oxygens and the cations. In the series studied the ratio O/Si is constant, therefore no change in the covalent cross-linking density occurs. Hence, in this case, any differences in T_{g} value must be attributed to changes of the numbers and strengths of cross-linking of the network modifier cations, having their own coordination and field strength, Z/r^2 (where Z is the charge and r the radius of the cation) [18,20]. Here, Ca²⁺ and La³⁺ have closely similar field strength values, $Z/r^2 =$ 2.04 \AA^{-2} , while small differences in the average coordination number in the glass can be expected as a consequence of the greater radius of the La^{3+} ion (1.21 Å instead of 0.99 Å). Fig. 4 shows no variations of T_{g} as La₂O₃ substitutes CaO. This trend of T_{g} curve is also consistent with previous results [4,15]: when the cations have similar Z/r^2 values, changes due to different coordination number occur when the amount substituted is sufficiently great. This was interpreted as the result of some competition between the two cations to achieve their preferred coordination. Therefore, the experimental results suggest that in the quenched glasses the lanthanum ion does not succeed in realizing its own average coordination. Whereas T_{g} depends on the structure of the glasses, T_p depends on the kinetic parameters of the devitrification process.

When considering the activation energy for crystal growth it is worth remembering that E_c is usually equal to the viscous flow activation energy E_{η} [16]. Because the structure rigidity does not vary and T_p shows little changes with the substitution, E_{η} and therefore E_c are expected to be constant; this is confirmed by our experimental results.

The *m* values appear to increase as the specific surface is increased (Fig. 11). Usually [21] the opposite result is obtained owing to the fact that the greater the specific surface the greater the tendency to devitrify by growth from surface nuclei so that *m* should be

progressively reduced to m=1. In the case of diopside glass [22] and glasses of composition CaO·SiO₂, 1.6CaO·0.4MgO·2SiO₂ and 1.4CaO·(0.6/3)Y₂O₃· 2SiO₂, it was found that nucleation preferentially occurs at the surfaces of the sample, but surface nuclei, formed in the glass transformation range, behave as bulk nuclei in finely powdered samples that efficiently sinter before devitrifying. This hypothesis appears to be effective in explaining the devitrification behavior of the studied glasses. In fact, in the case of coarsely powdered samples, that badly sinter, lower *m* values were obtained.

The evolution of XRD patterns shown in Fig. 5 is easily explained. The formation of a calcium poorer secondary phase, when CaO is substituted with La₂O₃ (2CaO·SiO₂ instead of 3CaO·SiO₂), can be attributed to the progressive reduction of the amount of calcium available. In the devitrified glass richest in La₂O₃, the majority phase is attributed to Ca₃La₆(SiO₄)₆ crystals. The growth of the α CaO·SiO₂ phase at lower temperatures than those shown in the phase diagram is easily explained considering that a sample subjected to a non-isothermal treatment, does not necessarily approaches equilibrium. Therefore, the α CaO·SiO₂ phase forms at lower temperature as its growth rate is high.

5. Conclusions

The glass transformation temperature, T_g , plots suggest that the substitution of La₂O₃ to CaO caused no variations of structure rigidity. In all devitrified samples, the lines of wollastonite (JCPDS card 27-88) and of a secondary calcium silicate phase appear. The XRD pattern of the free La₂O₃ sample also shows the lines of pseudo-wollastonite (JCPDS card 19/248) which in the CaO–SiO₂ phase diagram is reported to be stable only above 1125°C. It is relevant that, on the contrary, the devitrified product was obtained in the temperature range 900–1000°C. In the *x*=0.6 and *x*=1.0 devitrified samples, another crystalline phase is formed whose reflections were attributed to Ca₃La₆- $(SiO_4)_6$ crystals (JCPDS card 27/78). The activation energy for crystal growth E_c was 640 ± 25.1 kJ/mol regardless of composition. Devitrification involves a mechanism of surface nucleation, but, owing to softening and sintering, surface nuclei behave as bulk nuclei.

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